## BRIEF COMMUNICATIONS

## CONCERNING THE EFFECT OF INTERPHASE HEAT

## EXCHANGE ON THE LONGITUDINAL THERMAL

CONDUCTIVITY IN A GRANULAR LAYER
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Relations are derived for the temperature of elements in a granular layer and in a gas stream flowing against a uniform thermal flux, taking into account the interphase heat exchange.

In order to experimentally determine the longitudinal thermal conductivity in a granular layer ventilated by a gas stream, one usually applies the counterflow method to uniform thermal fluxes and gas streams [2,3] and, for this purpose, a flat heater is installed where the gas exits from the layer.

The steady-state heat exchange for this process configuration is described by the equations:

$$
\begin{gather*}
\lambda_{\mathrm{c}} \frac{\partial^{2} t}{\partial x^{2}}+c_{\mathrm{p}} G \frac{\partial t}{\partial x}=\alpha S\left(t-t_{\mathrm{c}}\right)  \tag{1}\\
\lambda_{0} \frac{\partial^{2} t_{\mathrm{c}}}{\partial x^{2}}=\alpha S\left(t_{\mathrm{c}}-t\right) . \tag{2}
\end{gather*}
$$

In dimensionless form we have

$$
\begin{gather*}
\frac{\partial^{2} \Theta}{\partial Y^{2}}+C \frac{\partial \Theta}{\partial Y}=C\left(\Theta-\Theta_{\mathrm{c}}\right) ;  \tag{3}\\
\frac{\partial^{2} \Theta_{c}}{\partial Y^{2}}=B\left(\Theta_{c}-\Theta\right) \tag{4}
\end{gather*}
$$

where

$$
\begin{gather*}
Y=\frac{\alpha S}{c_{\mathrm{p}} G} x=\frac{\mathrm{Nu}_{\mathrm{e}}}{\mathrm{Re}_{\mathrm{e}} \operatorname{Pr}} \cdot \frac{4}{d_{\mathrm{e}}} x ;  \tag{5}\\
B=\frac{\left(c_{\mathrm{p}} G\right)^{2}}{\alpha S} \cdot \frac{1}{\lambda_{0}}=\frac{(\operatorname{Re} \operatorname{Pr})^{2}}{\mathrm{Nue}^{2}} \cdot \frac{\varepsilon}{4 \bar{\lambda}_{0}} ;  \tag{6}\\
C=\frac{\left(c_{\mathrm{p}} G\right)^{2}}{\alpha S} \cdot \frac{1}{\lambda_{\mathrm{c}}}=B \frac{\lambda_{0}}{\lambda_{\mathrm{c}}} ;  \tag{7}\\
\Theta=\frac{t-t_{0}}{t_{\mathrm{c}, 1}-t_{0}} ; \quad \Theta_{c}=\frac{t_{\mathrm{c}}-t_{0}}{t_{\mathrm{c}, \mathrm{I}}-t_{0}} .
\end{gather*}
$$

The boundary conditions are

$$
\begin{equation*}
x=0, \quad \Theta_{c}=1 ; \quad x=\infty, \quad \Theta=\Theta_{c}=0 \tag{8}
\end{equation*}
$$

If the rate of interphase heat exchange is considerably higher than that of longitudinal heat transfer ( $\alpha S \geqslant c_{p} G$ ), then $\Theta$ and $\circlearrowleft_{C}$ were assumed close to one another so that Eq. (1) and (2) becomes

$$
\begin{equation*}
\lambda \frac{\partial^{2} \Theta}{\partial x^{2}}=-c_{\mathrm{p}} G \frac{\partial \Theta}{\partial x}, \tag{9}
\end{equation*}
$$

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Fig. 1. Longitudinal convective heat conductivity in a granular layer, and ratio of gas elements and layer elements temperatures when gas and heat are counterflowing with interphase heat exchange taken into account: 1) $\left.\bar{\lambda}_{0}=8,2\right) \bar{\lambda}_{0}=13$.
where $\lambda=\lambda_{0}+\lambda_{c}$. In this case

$$
\begin{align*}
& \Theta=\exp \left(-\frac{c_{\mathrm{p}} G}{\lambda} x\right)  \tag{10}\\
& m=\frac{\partial(\ln \Theta)}{\partial x}=-\frac{c_{\mathrm{p}} G}{\lambda} . \tag{11}
\end{align*}
$$

The magnitude of $\lambda$ was determined in [2, 3] according to Eq. (11). The solution to Eqs. (3) and (4) is then sought in the form:

$$
\begin{align*}
& \Theta_{\mathrm{c}}=\exp (-K Y)  \tag{12}\\
& \Theta=A \exp (-K Y) . \tag{13}
\end{align*}
$$

After inserting (12) and (13) into (3) and (4), we have

$$
\begin{gather*}
K^{2}=B\left(1-\frac{1}{1+K-\frac{K^{2}}{C}}\right)  \tag{14}\\
A=1-\frac{K^{2}}{B} \tag{15}
\end{gather*}
$$

From (12) and (13) we obtain

$$
\begin{equation*}
m=-K Y_{1} \tag{16}
\end{equation*}
$$

where $Y_{1}$ is taken from (5) at $x=1 \mathrm{~m}$. Taking into account (5) and (6), we rewrite expression (11) as

$$
\begin{equation*}
m=-B Y_{1} \frac{\lambda_{0}}{\lambda} \tag{17}
\end{equation*}
$$

Combining (16) and (17) will yield the ratio between the "apparent" value $\lambda_{c}$, app $=\lambda_{C}-\lambda_{0}$ found by test according to (10) and the real value $\lambda_{c}$ :

$$
\begin{equation*}
\mu \equiv \frac{\lambda_{\mathrm{c}, \mathrm{app}}}{\lambda_{\mathrm{c}}}=\left(\frac{B}{K}-1\right) \frac{\lambda_{0}}{\lambda_{\mathrm{c}}} . \tag{18}
\end{equation*}
$$

In accordance with the data in [1] on interphase heat exchange and longitudinal convective diffusion, the following values were used for calculating $\mu$ :

$$
\begin{gather*}
\mathrm{Re}_{\mathrm{e}}=1-100 ; \quad \mathrm{Nu}_{\mathrm{e}}=0.63 \operatorname{Re}_{\mathrm{e}}^{0,5 \mathrm{Pr}^{0,33}}  \tag{19}\\
\bar{\lambda}_{\mathrm{c}} \approx 0.28+0.5 \operatorname{Re}_{\mathrm{e}} \operatorname{Pr} \tag{20}
\end{gather*}
$$

The values of $B$ and $C$ were determined from Eqs. (6) and (7) with $\operatorname{Pr}=0.7$ and $\varepsilon=0.4$ for $\bar{\lambda}_{0}=8$ (glass) and 13 (steel), while $K$ was found from Eq. (14). The results of calculations within the range of $\operatorname{Re}_{\mathrm{e}}$ numbers encountered in tests [2] are shown in Fig. 1.

The conclusion drawn in [3] as to the difference between temperatures $\Theta$ and $\Theta_{c}$ being negligible is not accurate. The values of $\lambda_{\mathrm{c}}$ obtained in [2] are on the average $40 \%$ high. After appropriate corrections, they approach those calculated by Eq. (20). In Fig. 1 is also shown the ratio of temperatures $A=\Theta^{/\left(\Theta_{C}\right.}$ calculated by Eq. (15); its magnitude is slightly dependent on $\bar{\lambda}_{0}$.

| $c_{p}$ $d_{\text {e }}$ | specific heat of gas; equivalent diameter of granular layer; |
| :---: | :---: |
| G | mass rate of gas flow; |
| S | heat exchange surface per unit layer volume; |
| t | gas temperature; |
| $t_{c}$ | temperature of layer elements; |
| $\mathrm{t}_{0}$ | temperature of gas at entrance to layer; |
| $\mathrm{t}_{\mathrm{c}, 1}$ | temperature of layer elements at $\mathrm{x}=0$; |
| x | linear coordinate in the direction of heat flow through the layer; |
| $\alpha$ | coefficient of heat exchange between layer and gas elements; |
| $\lambda_{c}$ | longitudinal convective heat conductivity; |
| $\underline{\lambda}_{0}$ | thermal conductivity of unventilated layer; |
| $\bar{\lambda}=\lambda \sqrt{\lambda} g ;$ |  |
| $\lambda_{\mathrm{g}}$ | thermal conductivity of gas. |

## LITERATURE CITED

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