## CONCERNING THE EFFECT OF INTERPHASE HEAT EXCHANGE ON THE LONGITUDINAL THERMAL CONDUCTIVITY IN A GRANULAR LAYER

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UDC 536.244: 541.182.3

Relations are derived for the temperature of elements in a granular layer and in a gas stream flowing against a uniform thermal flux, taking into account the interphase heat exchange.

In order to experimentally determine the longitudinal thermal conductivity in a granular layer ventilated by a gas stream, one usually applies the counterflow method to uniform thermal fluxes and gas streams [2, 3] and, for this purpose, a flat heater is installed where the gas exits from the layer.

The steady-state heat exchange for this process configuration is described by the equations:

$$\lambda_{\rm c} \frac{\partial^2 t}{\partial x^2} + c_{\rm p} G \frac{\partial t}{\partial x} = \alpha S \left( t - t_{\rm c} \right); \tag{1}$$

$$\lambda_0 \frac{\partial^2 t_c}{\partial x^2} = \alpha S(t_c - t).$$
<sup>(2)</sup>

In dimensionless form we have

$$\frac{\partial^2 \Theta}{\partial Y^2} + C \frac{\partial \Theta}{\partial Y} = C \left( \Theta - \Theta_c \right); \tag{3}$$

$$\frac{\partial^2 \Theta_{\rm c}}{\partial Y^2} = B \left( \Theta_{\rm c} - \Theta \right),\tag{4}$$

where

$$Y = \frac{\alpha S}{c_n G} x = \frac{N u_e}{R e_o P r} \cdot \frac{4}{d_e} x;$$
(5)

$$B = \frac{(c_p G)^2}{\alpha S} \cdot \frac{1}{\lambda_0} = \frac{(\text{Re}_e \text{Pr})^2}{\text{Nue}} \cdot \frac{\varepsilon}{4\overline{\lambda_0}};$$
(6)

$$C = \frac{(c_{\rm p}G)^2}{\alpha S} \cdot \frac{1}{\lambda_{\rm c}} = B \frac{\lambda_0}{\lambda_{\rm c}};$$
<sup>(7)</sup>

$$\Theta = \frac{t - t_0}{t_{c,1} - t_0} ; \quad \Theta_c = \frac{t_c - t_0}{t_{c,1} - t_0} .$$

The boundary conditions are

$$x = 0, \quad \Theta_c = 1; \quad x = \infty, \quad \Theta = \Theta_c = 0.$$
 (8)

If the rate of interphase heat exchange is considerably higher than that of longitudinal heat transfer  $(\alpha S \gg c_{D}G)$ , then  $\Theta$  and  $\Theta_{C}$  were assumed close to one another so that Eq. (1) and (2) becomes

$$\lambda \frac{\partial^2 \Theta}{\partial x^2} = -c_p G \frac{\partial \Theta}{\partial x} , \qquad (9)$$

I. I. Polzunov Boiler and Turbine Institute, Moscow Division. Translated from Inzhenerno-Fizicheskii Zhurnal, Vol. 20, No. 2, pp. 344-346, February, 1971. Original article submitted January 21, 1970.

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Fig. 1. Longitudinal convective heat conductivity in a granular layer, and ratio of gas elements and layer elements temperatures when gas and heat are counterflowing with interphase heat exchange taken into account: 1)  $\overline{\lambda}_0 = 8$ , 2)  $\overline{\lambda}_0 = 13$ .

where  $\lambda = \lambda_0 + \lambda_c$ . In this case

$$\Theta = \exp\left(-\frac{c_{\rm p}G}{\lambda} x\right); \tag{10}$$

$$m \equiv \frac{\partial (\ln \Theta)}{\partial x} = -\frac{c_{\rm p}G}{\lambda}.$$
 (11)

The magnitude of  $\lambda$  was determined in [2, 3] according to Eq. (11). The solution to Eqs. (3) and (4) is then sought in the form:

$$\Theta_c = \exp\left(-KY\right);\tag{12}$$

$$\Theta = A \exp\left(-KY\right). \tag{13}$$

After inserting (12) and (13) into (3) and (4), we have

$$K^{2} = B\left(1 - \frac{1}{1 + K - \frac{K^{2}}{C}}\right);$$
(14)

$$A = 1 - \frac{K^2}{B} \,. \tag{15}$$

From (12) and (13) we obtain

$$m = -KY_1, \tag{16}$$

where  $Y_1$  is taken from (5) at x = 1 m. Taking into account (5) and (6), we rewrite expression (11) as

$$m = -BY_1 \frac{\lambda_0}{\lambda}.$$
 (17)

Combining (16) and (17) will yield the ratio between the "apparent" value  $\lambda_{c,app} = \lambda_c - \lambda_0$  found by test according to (10) and the real value  $\lambda_c$ :

$$\mu \equiv \frac{\lambda_{\rm c, app}}{\lambda_{\rm c}} = \left(\frac{B}{K} - 1\right) \frac{\lambda_0}{\lambda_{\rm c}}.$$
(18)

In accordance with the data in [1] on interphase heat exchange and longitudinal convective diffusion, the following values were used for calculating  $\mu$ :

$$Re_e = 1 - 100; Nu_e = 0.63 Re_e^{0.5} Pr^{0.33};$$
 (19)

$$\overline{\lambda_{c}} \approx 0.28 + 0.5 \text{Re}_{p} \text{Pr.}$$
(20)

The values of B and C were determined from Eqs. (6) and (7) with Pr = 0.7 and  $\varepsilon = 0.4$  for  $\overline{\lambda}_0 = 8$  (glass) and 13 (steel), while K was found from Eq. (14). The results of calculations within the range of Re<sub>e</sub> numbers encountered in tests [2] are shown in Fig. 1.

The conclusion drawn in [3] as to the difference between temperatures  $\Theta$  and  $\Theta_c$  being negligible is not accurate. The values of  $\lambda_c$  obtained in [2] are on the average 40% high. After appropriate corrections, they approach those calculated by Eq. (20). In Fig. 1 is also shown the ratio of temperatures  $A = \Theta / \Theta_c$  calculated by Eq. (15); its magnitude is slightly dependent on  $\overline{\lambda_0}$ .

c <sub>n</sub>	specific heat of gas;
de	equivalent diameter of granular layer;
G	mass rate of gas flow;
S	heat exchange surface per unit layer volume;
t	gas temperature;
te	temperature of layer elements;
t <sub>0</sub>	temperature of gas at entrance to layer;
te.1	temperature of layer elements at $x = 0$ ;
x	linear coordinate in the direction of heat flow through the layer;
α	coefficient of heat exchange between layer and gas elements;
λc	longitudinal convective heat conductivity;
λ <sub>0</sub>	thermal conductivity of unventilated layer;
$\overline{\lambda} = \lambda / \overline{\lambda}_{\rm cr};$	
λg	thermal conductivity of gas.
0	

## LITERATURE CITED

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